

ANALYSIS OF TM NONLINEAR OPTICAL WAVEGUIDE SENSORS

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Abstract

In this paper, we offer an analytical theory to study three-layer slab waveguide for sensor applications. Proposed sensor consists of a linear thin dielectric film that is surrounded by a nonlinear asymmetric environment. We consider polarized waves TM and nonlinear environments as of Kerr type. We also investigate sensor sensitivity for maximum mood, nonlinear sensors have a higher sensitivity than linear sensors. This sensitivity in nonlinear sensors can increase up to 10/072 % compared to linear sensors.

Keywords: Sensitivity, Thin dielectric film, Sensor fiber.

1. INTRODUCTION

In the last two decades, optical waveguides have been studied severely as sensor elements and have been taken place due to the growth of the technology and fiber optic [1-5]. These sensors have been studied for different purposes including humidity sensors, chemical sensors and biochemical sensors[6-12]. These types of sensors have important applications in determination of protein absorption based on dependency of bacteria recognition or cells living. Also, they can be used for detection of ultrathin biological molecules layers of thickness much smaller than guidance light wavelength and detection of harmful gases like methane, CO_2 , The privileges of these optical sensors are their immunity against electromagnetic interference, high sensitivity, small size and low price.

One of the important factors in the optical field is effective refractive index N which will be resulted by the $N = \frac{c}{v_{ph}}$ equation. In this equation C is the light velocity in vacuum and v_{ph} is the phase velocity. The main idea of planar dielectric waveguide sensor is measuring the changes in effective refractive index N due to changes in cover refractive index n_c . If a Light with the angle of α enter one of the ends of the waveguide in film part, it will be led within the film by total internal refraction which takes place at the boundaries of the film-cover and film-substrate. The mentioned light exits from the other end of the film in the angle of θ where the intensity is measured by a detector. The result of this devices will be the crude data which will be used in the drawing of the sensorgram chart [13-20].

Some of the injected light will be absorbed by the cover. Now if the cover was objected to an external evanescent field, the amount of the light absorption and phase displacement will change. These changes represent the density and the refractive index of the cover[21-23].

Now at this point an analysis will be performed to reach the best sensitivity by changing the properties of the system.

Afterwards the recommended sensor will be compared with the linear sensor of its own type.

2. THEORY

A schematic structure of the waveguide under consideration is shown in Figure 1. A guiding layer with permittivity ϵ_f and thickness h is coated onto a nonlinear substrate with permittivity ϵ_{nl3} . A nonlinear cladding layer with permittivity ϵ_{nl1} is coated onto the guiding layer. We will consider p-polarized waves that propagate in the x-direction (TM waves). The only nonvanishing components of the fields E and H are H_y , E_x , and E_z . Assuming the nonlinear dielectric functions to be of Kerr type, i.e., $\epsilon_{nl1} = \epsilon_c + \alpha_c |E|^2$ and $\epsilon_{nl3} = \epsilon_s + \alpha_s |E|^2$, where α_c and α_s are the nonlinear coefficients of the cladding and substrate, respectively and ϵ_c and ϵ_s are the linear parts of the permittivities. To solve the nonlinear wave equation for the magnetic field H , one can write ϵ_{nl1} and ϵ_{nl3} as :

$$\epsilon_{nl1} = \epsilon_c + \alpha'_c |H_{y1}|^2, \quad \epsilon_{nl3} = \epsilon_s + \alpha'_s |H_{y3}|^2 \quad (1)$$

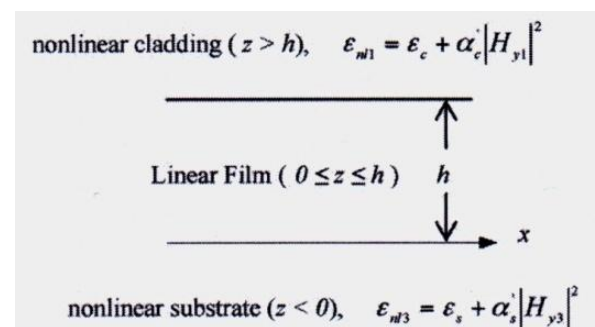


Fig- 1: Schematic structure of nonlinear slab waveguide sensor.

Where $\alpha'_c = \frac{\alpha_c}{\epsilon_c C^2 \epsilon_0^2}$ and $\alpha'_s = \frac{\alpha_s}{\epsilon_s C^2 \epsilon_0^2}$, C is the speed of light in vacuum, ϵ_0 is the free space permittivity and H_{y1} and H_{y3} are the TM fields in the cladding and substrate, respectively. After solving Maxwell's equations in the three layers of the

structure and matching the tangential magnetic and electric fields, the dispersion relations for $\alpha'_s > 0$ and $\alpha'_c > 0$ are given by:

$$k_0 q_f h - \text{rctan} \left(a \frac{X_c}{a_c} \tanh C_c \right) - \text{arctan} \left(\frac{X_s}{a_s} \tanh C_s \right) - m\pi = 0 \quad (2)$$

Where k_0 is the free space wave number, $q_f = \sqrt{\epsilon_f - N^2}$, N is the effective refractive index, $C_c = k_0 \sqrt{N^2 - \epsilon_c} (h - z_c)$, $C_s = k_0 \sqrt{N^2 - \epsilon_s} z_s$, z_c and z_s are constants related to the field distribution in the covering medium and substrate, respectively, $m=0,1,\dots$ is the mode order, a_s and a_c are two asymmetry parameters and X_s and X_c are two normalized variables given by:

$$a_s = \frac{\epsilon_s}{\epsilon_f}, \quad a_c = \frac{\epsilon_c}{\epsilon_f}, \quad X_s = \frac{\sqrt{N^2 - \epsilon_s}}{q_f}, \quad X_c = \frac{\sqrt{N^2 - \epsilon_c}}{q_f} \quad (3)$$

It is straightforward to show that X_s and X_c are interrelated by:

$$X_c^2 = \frac{(1 - a_c)(1 + X_s^2)}{(1 - a_s)} - 1 \quad (4)$$

The effective refractive index can be written in terms of a_s and X_s as

$$N = \sqrt{\epsilon_f} \sqrt{\frac{a_s + X_s^2}{1 + X_s^2}} \quad (5)$$

The sensitivity of the sensor S_h is defined as the rate of change of the effective refractive index under an index change of the cover. Differentiating Eq. (2), with respect to N and calculating S_h as $(\partial n_c / \partial N)^{-1}$ we obtain:

$$S_h = \frac{\sqrt{a_c} \sqrt{1 + X_c^2} B}{X_c \sqrt{a_c + X_c^2} D} \quad (6)$$

Where,

$$B = a_c H_c + a_c \tanh C_c + 2X_c^2 \tanh C_c (1 - a_c) / (1 + X_c^2) \quad (7)$$

$$D = (a_c^2 + X_c^2 \tanh^2 C_c) (A_{TM} + G_{sTM} + G_{cTM}) \quad (8)$$

$$H_c = k_0 (h - z_c) X_c \sqrt{\epsilon_f} \sqrt{\frac{1 - a_c}{1 + X_c^2}} (-\tanh^2 C_c) \quad (9)$$

$$G_{cTM} = \frac{a_c H_c + a_c \tanh C_c (1 + X_c^2)}{X_c (a_c^2 + X_c^2 \tanh^2 C_c)} \quad (10)$$

$$G_{sTM} = \frac{a_s H_s + a_s \tanh C_s (1 + X_s^2)}{X_s (a_s^2 + X_s^2 \tanh^2 C_s)} \quad (11)$$

$$H_s = k_0 (z_s) X_s \sqrt{\epsilon_f} \sqrt{\frac{1 - a_s}{1 + X_s^2}} (1 - \tanh^2 C_s) \quad (12)$$

$$A_{TM} = \text{arctan} \left(\frac{X_s}{a_s} \tanh C_s \right) + \text{arctan} \left(\frac{X_c}{a_c} \tanh C_c \right) + m\pi = 0 \quad (13)$$

The fraction of total power propagating in the covering medium is one of the most important quantities affecting the sensitivity of the sensor. For TM modes, the time averaged power flow in the x-direction per unit width in the y-direction can be expressed as

$$P = \frac{N k_0}{2\omega \epsilon_0 \epsilon_r} \int_{-\infty}^{+\infty} H_y^2 dz = P_s + P_f + P_c \quad (14)$$

Where P_s is the power flow in nonlinear substrate and P_f is the power flow in linear film.

The fraction of total power flowing in the nonlinear cladding is

$$\frac{P_c}{P_{total}} = \frac{\frac{X_c}{a_c \alpha_c} \sigma_c}{\frac{X_c}{a_c \alpha_c} \sigma_c + \frac{X_s^2 \text{sech}^2 C_s}{2\alpha_s'} r + \frac{X_s}{a_s \alpha_s} \sigma_s} \quad (15)$$

Where P_s is the power flow in nonlinear substrate and P_f is the power flow in linear film

Where,

$$r = k_0 q_f h x_+ + \frac{1}{2} \sin(2k_0 q_f h) x_- + X_s b / a_s \tanh C_s \quad (16)$$

$$\sigma_c = 1 - \tanh C_c \quad (17)$$

$$\sigma_s = 1 - \tanh C_s \quad (18)$$

$$b = 1 - \cos(2k_0 q_f h) \quad (19)$$

$$x_- = (1 - (X_s^2 / a_s^2)) \tanh^2 C_s \quad (20)$$

$$x_+ = (1 + (X_s^2 / a_s^2)) \tanh^2 C_s \quad (21)$$

3. REPRESENTATION AND DISCUSSION

To do the calculations, we have considered Si_3N_4 and GaN as the guiding layer. The refractive index of the mentioned materials is respectively $n_f = 2$ and $n_f = 2/34949$. The free space wavelength $\lambda = 1550$ nm, $\tanh C_c = 0.6$ and $\tanh C_s = 0.7$ is assumed. Due to the high sensitivity of the sensor, we considered fundamental mode $m=0$.

Covering and substratum medium are considered silica with different impurity percentages. It is necessary to mention that silica's linear and nonlinear dielectric constant is respectively $\epsilon = 2.13$ and $\alpha = 7.04 \times 10^{-20}$.

In figure 2, we have plotted a chart which represents the sensor sensitivity according the thickness of the wave guide film. The subjected sensor is of the symmetric type ($n_s > n_c$). This diagram was compared linear sensor of it own type (figure 3). As you can see in the figure 2, the sensitivity of the nonlinear type is both higher and reachable at a thinner thickness of the film.

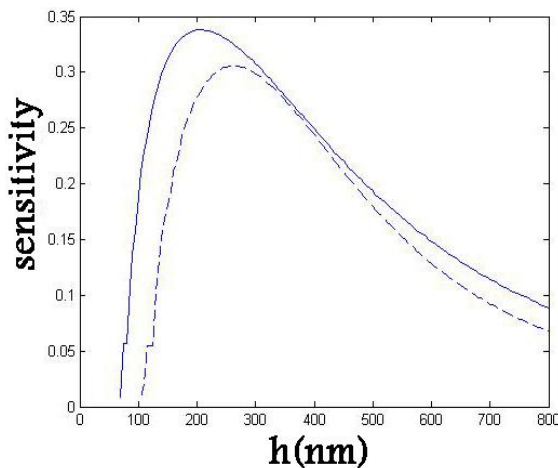


Fig- 2: Sensitivity versus the waveguide thickness h for $\epsilon_c = 2/13$ and $\epsilon_s = 2/343$ (Si_3N_4 , $n_f = 2$) for the proposed nonlinear sensor and linear sensor (dotted line).

The comparison which is done in Figures 3 is very important and interesting. According to the diagram, nonlinear sensors TM_0 have higher sensitivity in comparison with the TE_0 kinds but this maximum sensitivity occurs in a thicker thickness of the film and this is a very important point at the designing of the sensors. This diagram has been drawn for both environments of conduct layer i.e. Si_3N_4 and GaN which shows better sensitivity for GaN.

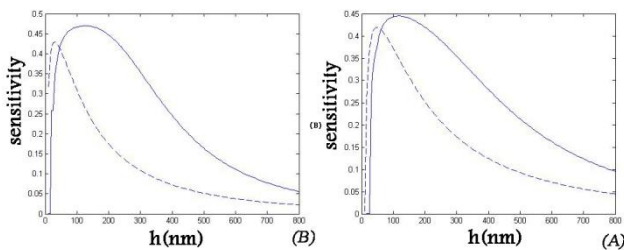


Fig- 3: Sensitivity as a function of waveguide film thickness for TM_0 (solid line) and TE_0 (dotted line) $\epsilon_c = 2/131917$, $\epsilon_s = 2/14917$ A- $n_f = 2$, (Si_3N_4) and B- $n_f = 2/34949$, (GaN).

In figure 4 the sensitivity is plotted as a function of effective refractive index N . The figure shows that the sensitivities are zero at $N = 1/57$ (this value of N occurs at cut-off since it is equal to n_s , the figure is plotted for $\epsilon_s = 2/4495$,

$n_s = \sqrt{\epsilon_s}$) and at $N = 2$ (this value of N is equal to the guiding layer refractive index). The sensitivities have their maxima somewhere between these two values near the cut-off and the values of the maxima of N , and also the possession of this maximum depends on the value of a_c . This behavior of the sensitivity with N is explained by the power considerations discussed below.

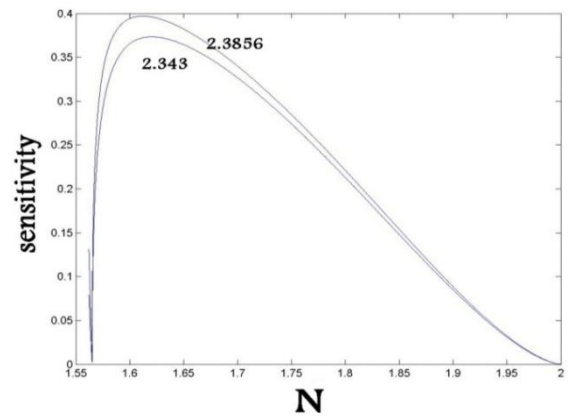


Fig-4 : Sensitivity versus the effective refractive index N for $\epsilon_c = 2.8856$, $\epsilon_c = 2.343$ and $\epsilon_s = 2/4495$, $n_f = 2$

Figure 5 verifies the close connection between the fraction of total power propagating in the nonlinear cladding medium (P_c/P_{total}) and the sensitivity of the sensor. In most cases, they may be regarded as nearly identical thus the enhancement of the fraction of power flowing in the cladding is essential for sensing applications.

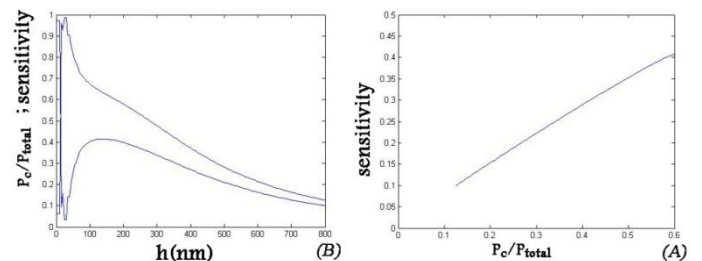


Fig- 5: Part A- sensitivity against the fraction of total power propagating in the cover medium, Part B-. the fraction of total power propagating in the cover medium as a function of the film thickness and the sensitivity as a function of the film thickness. ($\epsilon_s = 2/3856$, $\epsilon_c = 2/343$, $\alpha_c = 7/04 \times 10^{-21}$, $\alpha_s = 8/88 \times 10^{-21}$, $n_f = 2$)

4. CONCLUSIONS

In this work, optical sensor of nonlinear three-layer slab waveguide has been studied. The thickness of guiding layer is a critical factor for the sensitivity of the optical sensor.

As shown in figure 2, nonlinear sensors have higher sensitivity than linear sensors ones. Sensitivity of linear and nonlinear sensors (S_h) are respectively 0/3055 and 0/3378 that shows a 10/572% increase in the sensitivity for the nonlinear sensors type. Thickness of nonlinear and linear sensors (h) are respectively 200nm and 265 nm and that

nonlinear sensor shows decrease in thickness as much as 24/528%.

According to figure 3, nonlinear sensor sensitivity of TM_0 mode has higher amounts than TE_0 mode. Nonlinear sensor sensitivity in sensors with guiding layer GaN is higher than those with guiding layer Si_3N_4 . So sensors with guiding layer GaN and TM_0 mode are recommended

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