

SLOW STEADY MOTION OF A THERMO-VISCOUS FLUID BETWEEN TWO PARALLEL PLATES WITH CONSTANT PRESSURE AND TEMPERATURE GRADIENTS

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Abstract

In this paper, the slow steady motion of a second order thermo-viscous fluid between two parallel plates is examined. The closed form solutions of the velocity and temperature distributions are obtained when thermo-stress coefficient is far less compared to strain thermal conductivity coefficient and coefficient of cross viscosity for the following two cases: (i) when the upper plate is in relative motion and (ii) when the upper plate is thermally insulated. The heat transfer coefficient on the upper plate, The mean Bulk temperature and the transverse force perpendicular to the flow direction are also calculated. It is observed that forces are generated in transverse directions which are special feature of these types of fluids. The effect of various flow parameters on the flow field have been discussed with the help of graphical illustrations.

Keywords: Thermo-viscous fluids, Strain thermal conductivity coefficient and Thermo stress Coefficient.

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1. INTRODUCTION

The non-Newtonian nature of materials has been the subject of extensive study for over one and half centuries. It is only in last seven decades that serious attempts have been made to extend these investigations in the realm of non-linearity. The failure of the linear theories in predicting to a reasonable extent the mechanical behavior of materials such as liquid polymers, fluid plastic, the molten metal's etc subjected to stresses has been the motivating force behind study of the non-linear theories for material description. A non-linear generalization of the Voigt type materials was proposed by Rivlin [16] and Eringen [4]. Some of the non-linear theories proposed so far (listed in references) have not taken into account the strong dependence of visco-elastic behaviour upon thermal conduction i.e. interaction/interrelation between mechanical and non mechanical (such as thermal, chemical, electromagnetic etc.) effects even though the large amount data of experimental evidence indicate a strong dependence of visco-elastic nature of the fluid upon thermal behavior (Ferry [5]).

The development of non-linear theory reflecting the interaction/interrelation between thermal and viscous effects has been preliminarily studied by Koh and Eringen [9] and Coleman and Mizel [3]. A systematic rational approach for such a class of fluids has been developed by Green and Nagdhi [6]. In 1965 Kelly [10] examined some simple shear

flows of second order thermo-viscous fluids. Nageswara Rao and Pattabhi Ramacharyulu [14] later studied some steady state problems dealing with certain flows of thermo-viscous fluids. Some more problems K.Anuradha [1] and E.Nagaratnam [12] studied in plane, cylindrical and spherical geometries.

Flows of incompressible homogenous thermo-viscous fluids satisfy the following basic equations.

Equation of Continuity:

$$v_{i,i} = 0$$

Equation of Momentum:

$$\rho \left[\frac{\partial v_i}{\partial t} + v_k v_{i,k} \right] = \rho F_i + t_{ji,j}$$

Equation of Energy:

$$\rho c \dot{\theta} = t_{ij} d_{ij} - q_{i,i} + \rho \gamma$$

Where

$F_i = i^{th}$ Component of external force per unit mass

$c =$ Specific heat

γ = Thermal energy source per unit mass

and

$q_i = i^{th}$ Component of heat flux bivector $= \epsilon_{ijk} h_{jk} / 2$

Solving a specific boundary value problem would mean, finding the solution of these equations with appropriate boundary conditions such as the no slip condition (i.e. the velocity of fluid relative to the boundary is zero) and the prescription of the wall temperature. The later condition may be replaced by the prescription of heat flux on the boundary.

2. MATHEMATICAL FORMULATION AND SOLUTION OF THE PROBLEM

With reference to the Cartesian coordinates system $O(x, y, z)$ with origin on the lower plate, the X-axis in the direction of flow. The flow is characterized by the velocity field

$[u(y), 0, 0]$ and temperature by $\theta(y)$.

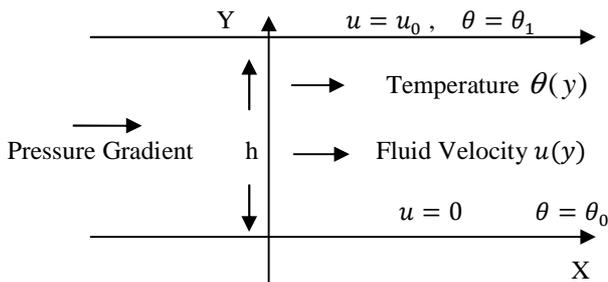


Fig 1: Flow Configuration

In the absence of any external force in the direction of flow, the equations of motion reduces to

$$0 = -\frac{\partial p}{\partial x} + \mu \frac{\partial^2 u}{\partial y^2} - \alpha_6 \frac{\partial \theta}{\partial x} \frac{\partial^2 \theta}{\partial y^2} \quad (1)$$

$$\mu_c \frac{\partial}{\partial y} \left(\frac{\partial u}{\partial y} \right)^2 + \rho F_y = 0 \quad (2)$$

$$\alpha_8 \frac{\partial}{\partial y} \left(\frac{\partial \theta}{\partial y} \frac{\partial u}{\partial y} \right) + \rho F_z = 0 \quad (3)$$

In the absence of any heat source, the energy equation reduces to

$$\rho c \left(u \frac{\partial \theta}{\partial x} \right) = \mu \left(\frac{\partial u}{\partial y} \right)^2 - \alpha_6 \frac{\partial \theta}{\partial x} \frac{\partial u}{\partial y} \frac{\partial \theta}{\partial y} + k \frac{\partial^2 \theta}{\partial y^2} + \beta_3 \frac{\partial \theta}{\partial x} \frac{\partial^2 u}{\partial y^2} \quad (4)$$

together with the boundary conditions:

$$u(0) = 0, \quad \theta(0) = \theta_0 \quad (5)$$

$$u(h) = u_0, \quad \theta(h) = \theta_1 \quad (6)$$

Introducing the following non-dimensional quantities,

$$y = hY, \quad u = (\mu/\rho h)U, \quad T = \frac{\theta - \theta_0}{\theta_1 - \theta_0}, \quad \frac{\partial \theta}{\partial x} = \frac{\theta_1 - \theta_0}{h} C_2$$

$$\text{and } -\frac{\partial p}{\partial x} = \frac{\mu^2}{\rho h^3} C_1$$

The equation (1) and (4) can be reduced to

$$0 = C_1 + \frac{d^2 U}{dY^2} - A_6 C_2 \frac{d^2 T}{dY^2} \quad (7)$$

and

$$U C_2 = A_1 \left[\left(\frac{dU}{dY} \right)^2 - A_6 C_2 \frac{dU}{dY} \frac{dT}{dY} \right] + \frac{1}{Pr} \frac{d^2 T}{dY^2} + B_3 C_2 \frac{d^2 U}{dY^2} \quad (8)$$

With

$$Pr = \frac{c\mu}{k} \text{ (Prandtl number)}, \quad S = \frac{\rho h v_0}{\mu}, \quad B_3 = \frac{\beta_3}{\rho h^2 c},$$

$$A_1 = \frac{\mu^2}{\rho h^2 c (\theta_1 - \theta_0)} \quad \text{and} \quad A_6 = \frac{\alpha_6 (\theta_1 - \theta_0)^2}{\mu^2}$$

where

C_1 is Non-Dimensional constant pressure gradient and C_2 is Non-Dimensional constant temperature gradient.

The boundary conditions are $U(0) = 0, \quad U(1) = U_0$

$$T(0) = 0, \quad T(1) = 1$$

3. CASE-I: WHEN UPPER PLATE IS IN RELATIVE MOTION

Assuming that the thermo – stress coefficient α_6 is far less when compared to strain thermal conductivity coefficient β_3 and coefficient of cross viscosity μ_c .

The equations (7) and (8) reduces to

$$0 = C_1 + \frac{d^2 U}{dY^2} \quad (9)$$

$$U C_2 = A_1 \left(\frac{dU}{dY} \right)^2 + \frac{1}{Pr} \frac{d^2 T}{dY^2} + B_3 C_2 \frac{d^2 U}{dY^2} \quad (10)$$

The boundary conditions reduces to

$$U(0) = 0, \quad T(0) = 0 \quad (11)$$

$$U(1) = U_0, \quad T(1) = 1 \quad (12)$$

i.e. the upper plate is moving with a given velocity and the two plates are maintained at different temperatures.

The equations (9) and (10) together with the boundary conditions (11) and (12) yield the velocity

$$U(Y) = \frac{C_1}{2} Y(1 - Y) + U_0 Y$$

And the temperature

$$T(Y) = Y + A_1 P_r \left[\frac{C_1^2}{24} (2Y^2 - 2Y + 1) + \frac{U_0^2}{2} - \frac{U_0 C_1}{6} (2 - 9Y) + \frac{B_3 C_2}{2} (-1 + Y + 2Y^2) \right]$$

The heat transfer coefficient i.e. Nusselt number “Nu” on the upper plate is

$$Nu = \left(\frac{dT}{dY} \right)_{Y=1} = -\frac{A_1 P_r}{24} [(C_1 - 2U_0)^2 + 8U_0^2] + \frac{C_2 P_r}{24} [C_1 + 8U_0 + 12B_3 C_2] + 1$$

Which depends on constant pressure gradient C_1 , the relative velocity of upper plate U_0 and B_3 the strain thermal conductivity coefficient, constant temperature gradient C_2 and the prandtl number p_r .

Also the mean Bulk temperature = $\frac{\int_0^1 uT \, dY}{\int_0^1 u \, dY}$

$$\left[\begin{array}{l} U_0 \left\{ \frac{1}{3} + A_1 P_r \left(\frac{C_1^2}{480} + \frac{U_0^2}{24} - \frac{U_0 C_1}{360} \right) \right\} \\ - C_2 P_r \left(\frac{C_1}{240} - \frac{U_0}{45} + \frac{B_3 C_2}{24} \right) \end{array} \right] \left/ \left(U_0 + \frac{C_1}{6} \right) \right.$$

$$+ C_1 \left\{ \frac{1}{24} + A_1 P_r \left(\frac{C_1^2}{2520} + \frac{U_0^2}{120} - \frac{U_0 C_1}{360} \right) \right\} \left/ \left(U_0 + \frac{C_1}{6} \right) \right.$$

$$- C_2 P_r \left(\frac{17C_1}{20160} - \frac{U_0}{240} + \frac{B_3 C_2}{240} \right) \left/ \left(U_0 + \frac{C_1}{6} \right) \right.$$

The forces generated in transverse direction are

$$\rho F_y = 2\mu_c \frac{\mu^2}{\rho^2 h^4} \left[U_0 + \frac{C_1}{2} (2Y - 1) \right] C_1$$

$$\rho F_z = \frac{\alpha_8 \mu (\theta_1 - \theta_0)}{h^4 \rho}$$

$$\left[\begin{array}{l} c_1 \left\{ \begin{array}{l} 1 + A_1 p_r \left\{ \frac{C_1^2}{6} (1 + 6Y^2 - 8Y^3) \right\} \\ + \frac{U_0 C_1}{3} (2 - 9Y) \end{array} \right\} \\ C_2 p_r \left\{ \begin{array}{l} -\frac{C_1}{24} (1 + 6Y - 24Y^2 + 16Y^3) \\ + \frac{U_0}{6} (1 + 3Y - 9Y^2) \\ - \frac{B_3 C_2}{2} (-1 + Y + 2Y^2) \end{array} \right\} \\ -U_0 p_r \left\{ \begin{array}{l} \frac{C_1^2}{4} (1 - 4Y + 4Y^2) - U_0^2 \\ - 3U_0 C_1 (-1 + 2Y) \\ + C_2 \left\{ \frac{C_1}{2} (Y - Y^2) - U_0 Y \right\} \\ + B_3 C_2 Y \end{array} \right\} \end{array} \right]$$

It is observed that these forces generated in transverse directions depends on the cross viscosity μ_c and on α_8 the thermo stress viscosity.

4. CASE-II: WHEN UPPER PLATE IS INSULATED

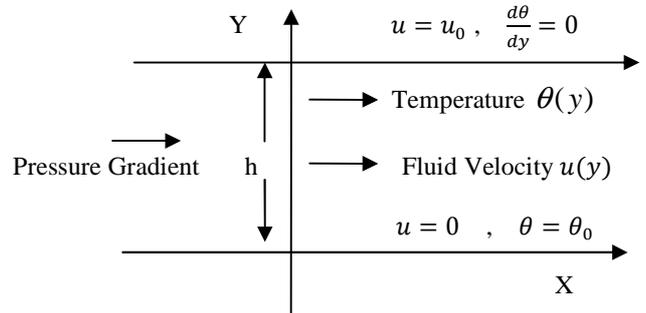


Fig.2: Flow Configuration

The equations of motion and energy reduces to

$$0 = c_1 + \frac{d^2 u}{dY^2} - A_6 C_2 \frac{d^2 T}{dY^2} \tag{13}$$

$$U C_2 = A_1 \left[\left(\frac{dU}{dY} \right)^2 - A_6 C_2 \frac{dU}{dY} \frac{dT}{dY} \right] + \frac{1}{p_r} \frac{d^2 T}{dY^2} + B_3 C_2 \frac{d^2 U}{dY^2} \tag{14}$$

Together with the boundary conditions:

$$U(0) = 0 \quad , \quad T(0) = 0 \tag{15}$$

$$U(1) = U_0 \quad , \quad \left(\frac{dT}{dY} \right)_{Y=1} = 0 \tag{16}$$

The velocity field is same as in Newtonian case and is given by

$$U(Y) = \frac{C_1}{2} Y(1 - Y) + U_0 Y$$

And the temperature field is given by

$$T(Y) = p_r C_2 \left[\frac{C_1}{24} (-Y^3 + 2Y^2 - 2) + \frac{U_0}{6} (Y^2 - 3) \right] Y + A_1 p_r \left[\frac{C_1^2}{24} (2 - 3Y + 4Y^2 - 2Y^3) - \frac{C_1 U_0}{6} (3 - 2YY + U_0 22(2 - Y)Y + B_3 C_1 C_2 p_r 2 (Y - 2)Y \right]$$

The forces generated in transverse direction are

$$\rho F_y = 2\mu_c \frac{\mu^2}{\rho^2 h^4} \left[U_0 + \frac{C_1}{2} (2Y - 1) \right] C_1$$

$\rho F_z =$

$$C_1 \alpha_8 \mu \theta_0 \left[\begin{array}{l} p_r C_2 \left\{ \frac{C_1}{12} (1 - 3Y + 6Y^2 - 4Y^3) + \frac{U_0}{2} (1 + Y - 3Y^2) \right\} \\ + p_r A_1 \left[\frac{C_1^2}{24} (32Y^3 - 48Y^2 + 24Y - 5) - C_1 U_0 (-5Y^2 + 5Y - 1) + \frac{U_0^2}{2} (4Y - 3) \right] \\ + \frac{p_r B_3 C_1 C_2}{2} (3 - 4Y) \end{array} \right] + U_0 \left[\begin{array}{l} p_r C_2 \left\{ \frac{C_1}{12} (-Y^2 + Y) + U_0 Y \right\} \\ - p_r A_1 \left[\frac{C_1^2}{24} (4Y^2 - 4Y + 1) + C_1 U_0 (1 - 2Y) + U_0^2 \right] \\ + p_r B_3 C_1 C_2 \end{array} \right]$$

Which depends on α_8 , B_3 , pressure gradient, given velocity of upper plate and the constant temperature gradient.

5. RESULTS AND DISCUSSION

Numerical estimates of the velocity and temperature fields was carried for different values of $U_0=(0, 1)$ by taking $C_1 = 1, A_1 = 1, p_r = 1$ and these are illustrated graphically.

When the upper plate is fixed, thermally insulated or not, the velocity is parabolic in general. When the upper plate is in

relative motion with a given velocity, the velocity of the fluid is steadily increased so as to attain the velocity of the upper plate, this is observed from the Fig 3.

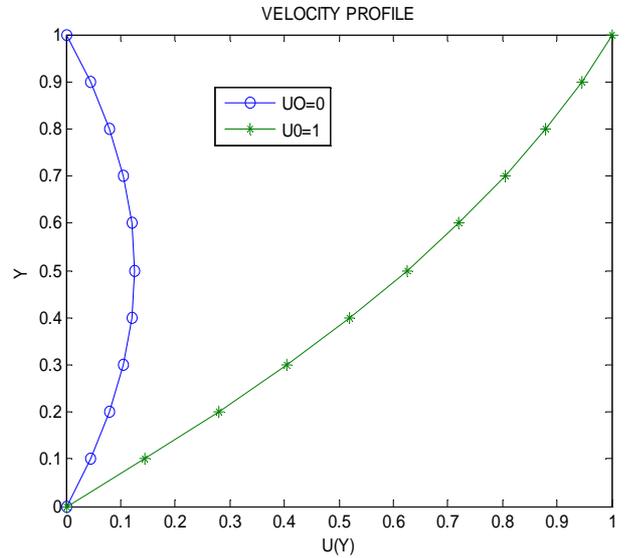


Fig 3: Velocity Profile

The temperature distributions for different values of B3 are illustrated graphically in Fig. 4, 5 and 6

When the upper plate is not thermally insulated temperature increases gradually to attain the temperature of the upper plate and when it is thermally insulated temperature decreases.

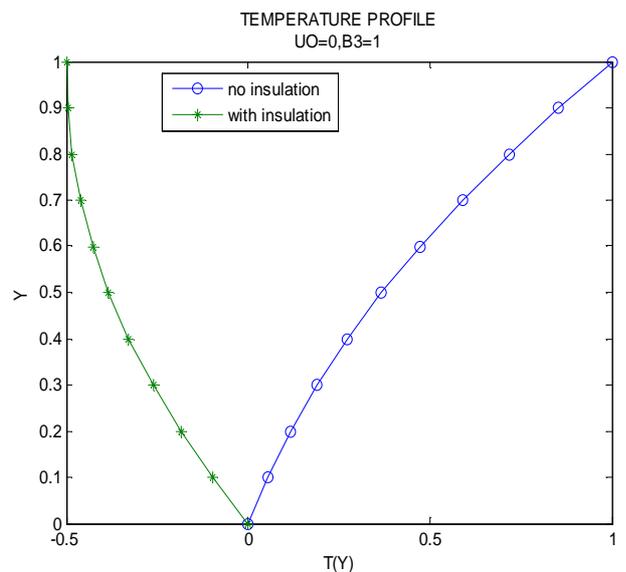


Fig 4: Temperature Profile

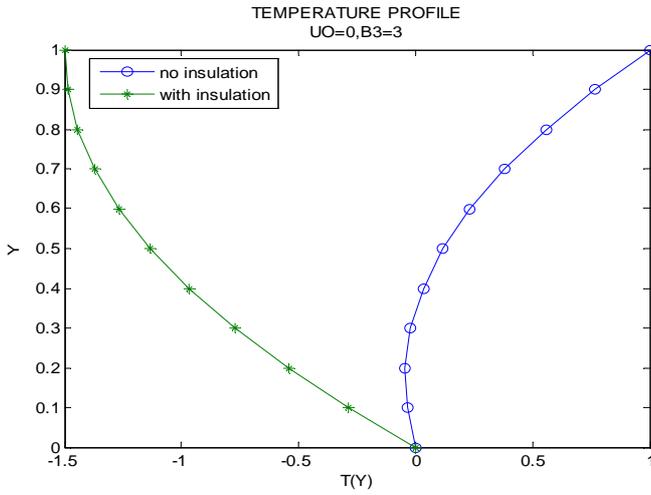


Fig.5: Temperature Profile

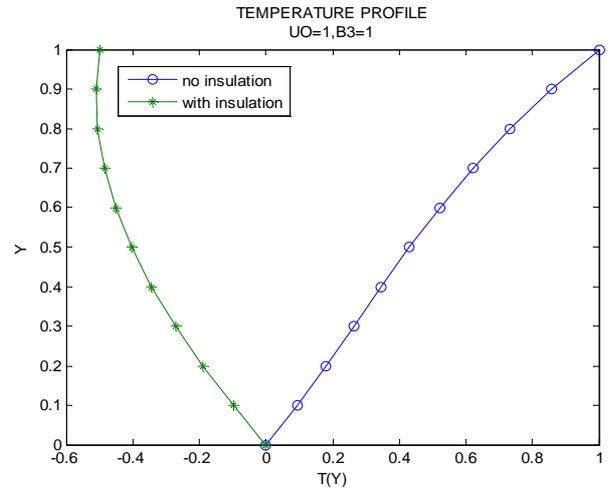


Fig.7: Temperature Profile

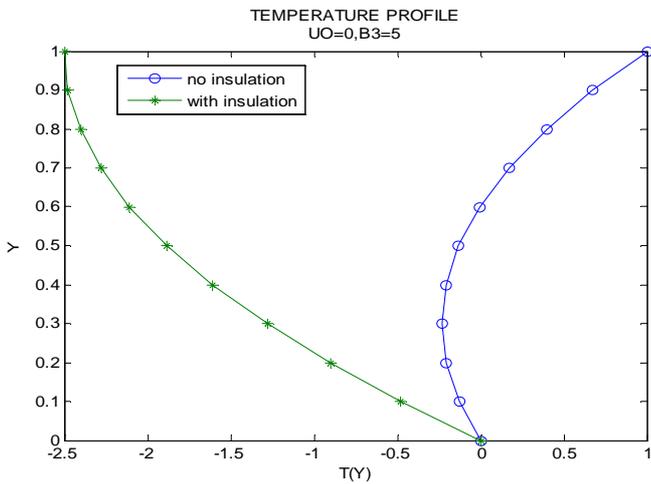


Fig.6: Temperature Profile

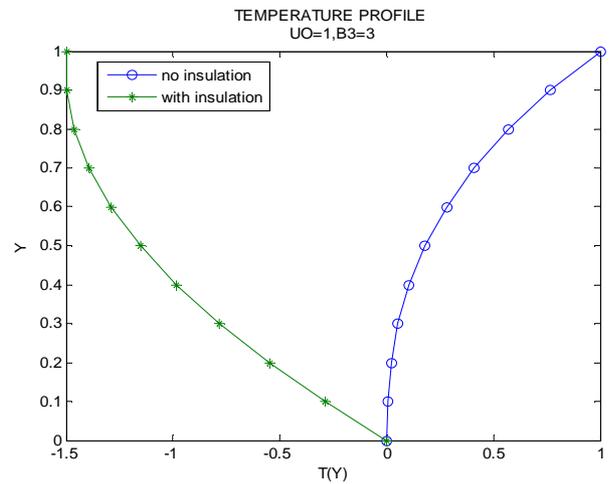


Fig.8: Temperature Profile

When the upper plate is in relative motion, temperature distributions for different values of $B3$ are illustrated graphically in the figures 7, 8 and 9.

When the upper plate is not thermally insulated temperature increases gradually to attain the temperature of the upper plate and when it is thermally insulated temperature decreases.

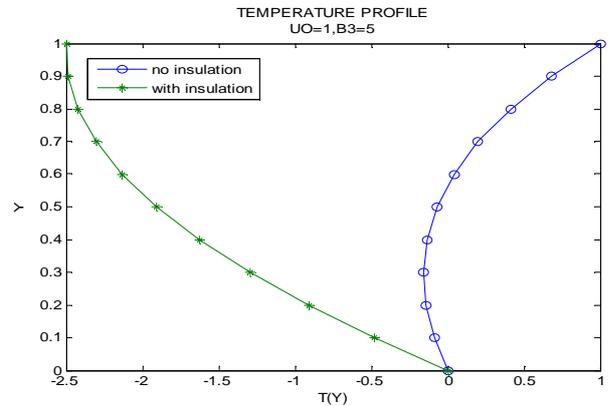


Fig.9: Temperature Profile

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